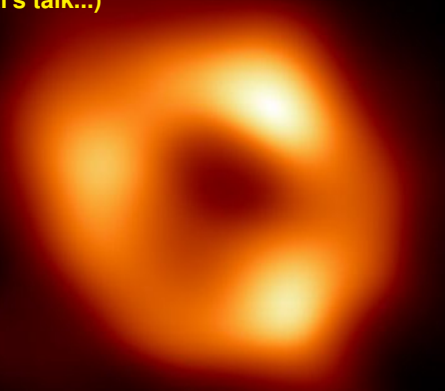


# Red giant - jet interaction in galactic nuclei

Hydrodynamical simulations of repetitive stellar passages

(following Michal's talk...)

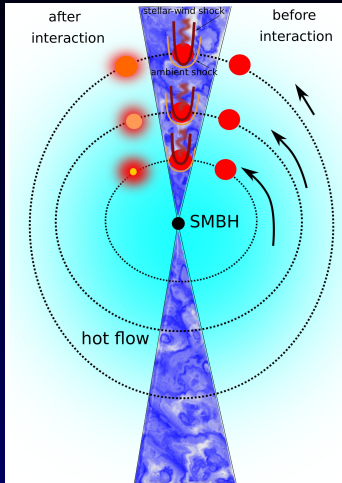


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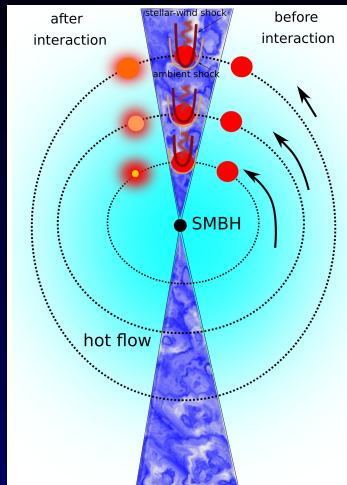
**CPB Meeting**, Brno, June 1, 2022

**Galactic center** - the inner  $\sim 1$  pc is a region of mutual interactions of stars, gas and dust within the gravitational potential of the SMBH



- illustration of the jet - red giant interaction
- at lower  $z$  this is expected to be stronger

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### Ambient medium:

The  $\rho$  and  $T$  profiles of the ambient plasma - power-law functions

$$n_a \approx n_B \left( \frac{r}{r_B} \right)^{-1}, \quad (1)$$

$$T_a \approx T_B \left( \frac{r}{r_B} \right)^{-1}, \quad (2)$$

where  $n_B = 26 \text{ cm}^{-3}$ , and  $T_B = 1.5 \times 10^7 \text{ K}$  are the number density and the temperature at the Bondi radius

$$r_B = \frac{2GM_\bullet}{c_s^2} \sim 0.21 \left( \frac{T_B}{10^7 \text{ K}} \right)^{-1} \text{ pc}, \quad (3)$$

where  $M_\bullet = 4 \times 10^6 M_\odot$

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### Jet structure:

We assume that the jet plasma is matter-dominated, consisting of electrons and protons. The jet exerts the pressure on the passing star mainly in the form of the bulk motion of the jet plasma at the velocity of  $v_j$ , which results in the ram pressure of  $P_j = \Gamma \rho_j v_j^2$ , where  $\Gamma$  is the Lorentz factor and  $\rho_j$  is the mass density inside the jet. The **number density** inside the hadronic jet can then be estimated as (Zajaček et al., 2020),

$$n_j = \frac{L_j}{\mu m_H (\Gamma - 1) c^2 v_j \pi Z^2 \tan^2 \theta} \\ \simeq 53 \left( \frac{L_j}{10^{42} \text{ erg s}^{-1}} \right) \left( \frac{z}{0.01 \text{ pc}} \right)^{-2} \text{ cm}^{-3}, \quad (4)$$

which gives the **mass density**  $\approx 10^{-18} \text{ g cm}^{-3}$  at  $10^{-3} \text{ pc}$ .

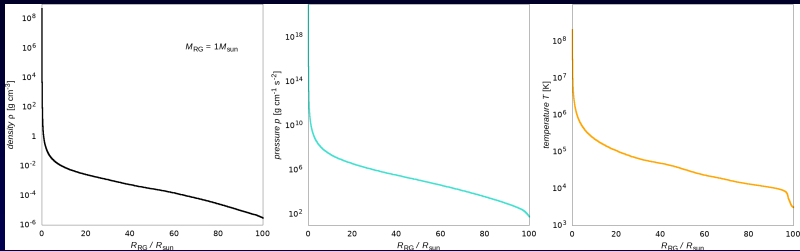
The **jet temperature** is assumed to be  $T_j = 10^{10} \text{ K}$  (Bosch-Ramon et al., 2012)

We assume the **jet luminosity**  $L_j = 10^{42} \text{ erg s}^{-1}$ , the **jet velocity**  $v_j = 0.3 c$ , and the **jet opening half-angle**  $10^\circ$

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### Red giant model:

We model the red-giant as a star with mass  $M_{\text{RG}} = 1 M_{\odot}$ , and the radius  $R_{\text{RG}} = 100 R_{\odot}$ . The initial profiles of density, pressure, and temperature are calculated using the **stellar evolution code MESA** (e.g., Paxton et al., 2010).



We select **sufficiently higher initial mass** of the star to obtain  $1 M_{\odot}$  and  $100 R_{\odot}$  RGB star before the He-flash

We remap the MESA **density, pressure, and temperature** profiles to our computational grid, using its refined structure towards the stellar center

## Global structure of the own hydrodynamic (MHD) code

(cf. Kurfürst & Krtička 2014, 2018; Kurfürst et al., 2017, 2019, 2020)

### Conservative equations of ideal MHD:

$$\partial_t \rho + \vec{\nabla} \cdot (\rho \vec{v}) = 0, \quad (5)$$

$$\partial_t (\rho \vec{v}) + \vec{\nabla} \cdot (\rho \vec{v} \vec{v} + \mathcal{P}) = (8\pi)^{-1} [2(\vec{B} \cdot \vec{\nabla}) \vec{B} - \vec{\nabla} B^2] + \rho \vec{g}, \quad (6)$$

$$\partial_t E + \vec{\nabla} \cdot [(E + \mathcal{P}) \cdot \vec{v}] = (8\pi)^{-1} \left\{ \vec{\nabla} \cdot [2(\vec{B} \cdot \vec{v}) \vec{B} - B^2 \vec{v}] \right\} + \rho \vec{g} \cdot \vec{v}, \quad (7)$$

$$\partial_t \vec{B} + \vec{B} \vec{\nabla} \cdot \vec{v} + (\vec{v} \cdot \vec{\nabla}) \vec{B} - (\vec{B} \cdot \vec{\nabla}) \vec{v} = \vec{0}, \quad (8)$$

- where  $\mathcal{P}$  is the pressure tensor (including shear terms),  $\vec{g} = \vec{g}_{\text{grav}} + \vec{g}_{\text{rot}} + \vec{g}_{\text{rad}}$ , and  $E = E_{\text{int}} + E_{\text{kin}} + E_{\text{mag}}$
- The scalar thermal pressure  $\rho$  follows the ideal MHD EOS:

$$\rho = (\gamma - 1) \left[ E - \rho v^2 / 2 - B^2 / (8\pi) \right] \quad (9)$$

- All the equations are complemented with the divergence-free constraint:

$$\vec{\nabla} \cdot \vec{B} = 0 \quad (10)$$

(We currently involve only the hydrodynamic part for the simulations!)

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(cf. Kurfürst & Krtička 2014, 2018; Kurfürst et al., 2017, 2019, 2020)

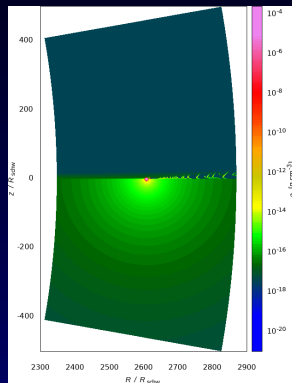
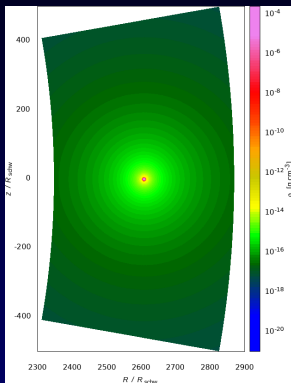
### Two types of hydro-solvers:

- operator-split (**HLLC**) finite volume Eulerian algorithm on staggered mesh (Stone & Norman 1992)
- unsplit Eulerian **Roe solver** (Roe 1981; Toro 1999) for strong shocks
- **MHD solver** for both types; for the Roe solver only in Cartesian form
- **all basic geometries** (Cartesian 3D, cylindrical 2.5D, spherical 2.5D) plus one non-orthogonal for “flaring” disks (Kurfürst & Krtička 2018)
- **Navier-Stokes viscosity** solver in all the geometries
- **static mesh refinement** (in this simulation 2700 / 3600 grid cells)
- full implementation of **MPI** for parallelization

Currently is being upgraded (among other purposes) for the 2D analogy of the SN explosion code SNEC

## Snapshots of the density

- orbital radius is 0.001 pc
- initial ambient stellar wind corresponds to  $\dot{M}_{\text{RG}} = 10^{-6} M_{\odot} \text{ yr}^{-1}$
- wind expansion velocity is  $15 \text{ km s}^{-1}$
- BCs are inflow at left and top, outflow at right at bottom

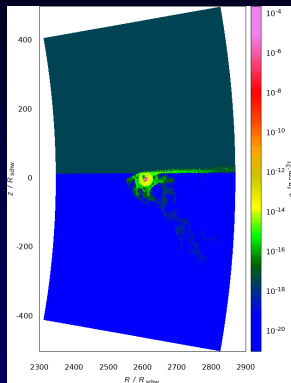
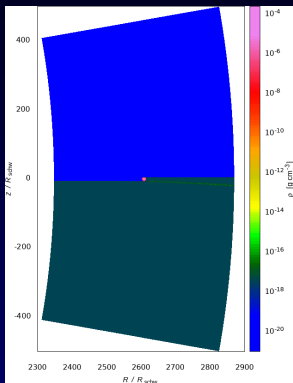


- **Left panel:** start of the simulation at  $t = 0$
- **Right panel:** first entry to the jet at  $t \approx 15$  d



## Snapshots of the density

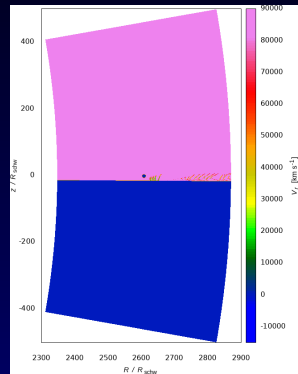
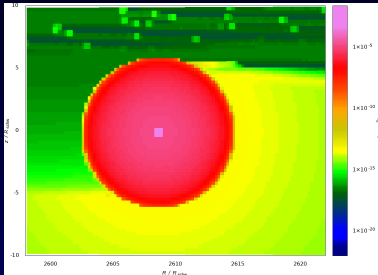
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- **Left panel:** first exit of the jet at  $t \approx 45$  d
- **Right panel:** second entry to the jet at  $t \approx 285$  d

## Snapshot of the density and velocity

- orbital radius is 0.001 pc
- initial ambient stellar wind corresponds to  $\dot{M}_{\text{RG}} = 10^{-6} M_{\odot} \text{yr}^{-1}$
- wind expansion velocity is  $15 \text{ km s}^{-1}$
- BCs are inflow at left and top, outflow at right at bottom



- **Left panel:** snapshot of the stellar region density at first jet entry
- **Right panel:** snapshot of the velocity at first jet entry (a bit boring)

## ... the videos of the first phase of the process

The simulation is calculated in the reference frame connected with the star

- ▶ Aligned dipole: [jet\\_cross\\_0.001pc\\_global\\_density.mp4](#)
  
  - ▶ Aligned quadrupole: [jet\\_cross\\_0.001pc\\_detail\\_density\\_phase1.mp4](#)
- Density integrations after first jet crosses reveal the stellar density loss  $\sim 0.0005$  per cross, however, that may be still uncertain - let's wait for a complete simulation including several hundreds or thousands passages



**Thank you!**